Partially dissipative systems: hypocoercivity and hyperbolisation

Timothée Crin-Barat

Friedrich-Alexander-Universität (FAU)

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1. **First part:** Stability of partially dissipative hyperbolic systems

2. **Second part:** Hyperbolisation via partial dissipation

Hypocoercivity for hyperbolic systems

We consider *n*-component hyperbolic systems of the form:

$$\begin{cases} \partial_t U + \sum_{j=1}^{d} A^j(U)\partial_{x_j} U + BU = 0, \\ U_0(x,t) = U_0(x), \end{cases} \tag{1}$$

such that $t > 0$, $x \in \mathbb{R}^d$, the matrices valued maps A are symmetric and B is a positive $n \times n$ matrix.

Three scenarios:

- When $B = 0$, for small and smooth data → local-in-time solutions (Kato, Majda, Serre) that may develop singularities (shock waves) in finite time (Dafermos, Lax).
- When $\mathrm{rank}(B) = n$, existence of global-in-time solutions (Li) that are exponentially damped.
- What can we say about the long-time behaviour in the partially dissipative setting: $0 < \mathrm{rank}(B) < n$?

We focus on the one-dimensional systems of the form

$$\partial_t U + A\partial_x U + BU = 0, \tag{2}$$

where A is symmetric, B is *partially dissipative*: $\mathrm{rank}(B) = n_2 < n$ where $n_1 + n_2 = n$ and

$$B = \begin{pmatrix} 0 & 0 \\ 0 & D \end{pmatrix} \quad \text{with } D > 0.$$

Additionally, we assume either one of these conditions:

1. D is symmetric.
2. **Or** D is strongly dissipative: for every $X \in \mathbb{R}^{n_2}$, there exists $\kappa > 0$ such that
$$\langle DX, X\rangle \geq \kappa \|X\|^2.$$

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Decomposing $U = (U_1, U_2)$, we have

$$\begin{cases} \partial_t U_1 + A_{1,1}\partial_x U_1 + A_{1,2}\partial_x U_2 = 0, \\ \partial_t U_2 + A_{2,1}\partial_x U_1 + A_{2,2}\partial_x U_2 = -DU_2, \end{cases} \quad \text{where } A = \begin{pmatrix} A_{1,1} & A_{1,2} \\ A_{2,1} & A_{2,2} \end{pmatrix}.$$

Applications

Example of applications: • The compressible Euler equations with damping:

$$\begin{cases} \partial_t \rho + \partial_x(\rho u) = 0, \\ \partial_t(\rho u) + \partial_x(\rho u^2) + \partial_x P(\rho) + \rho u = 0, \end{cases}$$

For the pressure law $P(\rho) = A\rho^\gamma$, with $A > 0$ and $\gamma > 1$, we can rewrite System (5) into the symmetric form:

$$\begin{cases} \partial_t c + u\partial_x c + \dfrac{\gamma - 1}{2} c\partial_x u = 0, \\ \partial_t u + u\partial_x u + \dfrac{\gamma - 1}{2} c\partial_x c = -u, \end{cases} \tag{3}$$

where $c = \sqrt{\frac{\partial P(\rho)}{\partial \rho}}$ corresponds to the sound speed.

• *Partial dissipation* occurs in many compressible models including dissipation: Compressible Navier-Stokes equations, Chemotaxis systems, Timoshenko systems, Discrete BGK, Euler-Maxwell equations, etc.

Time-decay rates in the linear setting.

Context

Goal: establish time-decay rates for

$$\partial_t U + A\partial_x U + BU = 0.$$

First difficulty: partial dissipation leads to an obvious lack of coercivity:

$$\frac{1}{2}\frac{d}{dt}\|(U_1,U_2)(t)\|_{L^2}^2 + \kappa\|U_2(t)\|_{L^2}^2 \leq 0, \tag{4}$$

→ no time-decay information on U_1.

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Inspiration to tackle this issue: Theories of hypoellipticity (Hörmander) and hypocoercivity (Villani):

"There might be regularizing/stabilizing mechanisms *hidden* in the interactions between the hyperbolic part A and the dissipative matrix B."

→ Let's see what this means in the context of ODEs.

ODE toy-model

Consider the ODE

$$\partial_t U + AU + BU = 0 \tag{5}$$

such that A is skew-symmetric and B positive symmetric (rank(B)$< n$).

If (A, B) satisfies the Kalman rank condition:

$$\mathrm{rank}(B, BA, BA^2, \dots, BA^{n-1}) = n \tag{K}$$

then the solutions of (5) satisfy

$$\|U(t)\|_{L^2} \leq C\|U_0\|_{L^2} e^{-\lambda t}.$$

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Sketch of proof: Again, since A is skew-symmetric, we have

$$\frac{1}{2}\frac{d}{dt}\|U(t)\|_{L^2}^2 + \kappa\|U_2(t)\|_{L^2}^2 \leq 0. \tag{6}$$

And, using the interactions between A and B,

$$\frac{d}{dt}\left(\sum_{k=1}^{n-1} < BA^{k-1}U, BA^k U >\right) + \sum_{k=1}^{n-1} \|BA^k U(t)\|_{L^2}^2 \leq C\|U_2(t)\|_{L^2}^2 + \dots$$

Under the Kalman rank condition, we have

$$\sum_{k=0}^{n-1} \|BA^k U(t)\|_{L^2}^2 \sim \|U(t)\|_{L^2}^2.$$

Therefore, the following functional is a Lyapunov functional

$$\mathcal{L}(t) = \|U(t)\|_{L^2}^2 + \eta \left(\sum_{k=1}^{n-1} < BA^{k-1}U, BA^k U >_{L^2} \right)$$

verifying

$$\frac{d}{dt}\mathcal{L}(t) + \|U_2(t)\|_{L^2}^2 + \eta\|U(t)\|_{L^2}^2 \leq \eta\|U_2(t)\|_{L^2}^2.$$

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$$\frac{d}{dt}\mathcal{L}(t) + \|U_2(t)\|_{L^2}^2 + \eta \|U(t)\|_{L^2}^2 \leq \eta \|U_2(t)\|_{L^2}^2.$$

For η small enough, we have

$$\mathcal{L}(t) \sim \|U(t)\|_{L^2}^2$$

and thus

$$\frac{d}{dt}\mathcal{L}(t) + \eta \mathcal{L}(t) \leq 0. \quad \square$$

Partially dissipative hyperbolic systems

- In the hyperbolic setting, the idea is essentially the same.

Main difficulty: The operators $A\partial_x$ and B are of a different order.

$\rightarrow$ Need to find a way to make them communicate as in the ODE setting.

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Two approaches:

- Fourier-based approach.

Roughly, one can proceed as in the ODE setting by adding frequency weights to the Lyapunov functional.

- Time-weighted Fourier-free approach.

→ Not optimal results but a broader range of applications (e.g. numerics, bounded domains, nonlinear dissipation).

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Toy-model analysis

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Let us look at the damped *p*-system:

$$\begin{cases} \partial_t \rho + \partial_x u = 0, \\ \partial_t u + \partial_x \rho + u = 0. \end{cases} \tag{7}$$

Again, standard energy estimates lead to

$$\frac{1}{2}\frac{d}{dt}\|(\rho,u)(t)\|_{L^2}^2 + \|u(t)\|_{L^2}^2 = 0.$$

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To overcome the lack of coercivity, we consider the Lyapunov functional:

$$\mathcal{L}_1(t) = \|(\rho, u, \partial_x \rho, \partial_x u)(t)\|_{L^2}^2 + \frac{1}{2}\int_{\mathbb{R}} u \partial_x \rho \, dx. \tag{8}$$

Notice that

$$\mathcal{L}_1(t) \sim \|(\rho, u, \partial_x \rho, \partial_x u)(t)\|_{L^2}^2.$$

Differentiating-in-time (8), we get

$$\frac{d}{dt}\mathcal{L}_1(t) + \|(u, \partial_x u)(t)\|_{L^2}^2 + \|\partial_x \rho(t)\|_{L^2}^2 \leq 0, \tag{9}$$

Detailed computations

$$\begin{cases} \partial_t \rho + \partial_x u = 0, \\ \partial_t u + \partial_x \rho + u = 0. \end{cases}$$

Standard H^1 estimates:

$$\frac{d}{dt} \|(\rho, \partial_x \rho, u, \partial_x u)\|_{L^2}^2 + \|(u, \partial_x u)\|_{L^2}^2 = 0$$

Cross estimates:

$$\frac{d}{dt} \int_{\mathbb{R}} u \partial_x \rho \, dx + \|\partial_x \rho\|_{L^2}^2 = \|\partial_x u\|_{L^2}^2 + \int_{\mathbb{R}} u \partial_x \rho.$$

Using Young inequality and gathering the estimates, we get

$$\frac{d}{dt} \mathcal{L}_1(t) + \|(u, \partial_x u)(t)\|_{L^2}^2 + \|\partial_x \rho(t)\|_{L^2}^2 \leq 0, \tag{10}$$

Second difficulty: how to get decay estimates from here?

Fourier heuristics

We have

$$\frac{d}{dt}\mathcal{L}_1(t) + \|(u, \partial_x u)(t)\|_{L^2}^2 + \|\partial_x \rho(t)\|_{L^2}^2 \leq 0. \tag{11}$$

Heuristically, in the frequency world, it reads

$$\frac{d}{dt}\mathcal{L}_1(t) + \|\min(1, \xi)(\widehat{u}, \widehat{\rho})\|_{L^2}^2 \leq 0. \tag{12}$$

From which it is easy to obtain

- A heat behavior for low frequencies,
- Exponential decay for high frequencies:

$$\|(\rho, u)^{\ell}(t)\|_{L^\infty} \leq Ct^{-1/2}\|(\rho_0, u_0)\|_{L^1}, \tag{13}$$

$$\|(\rho, u)^{h}(t)\|_{L^2} \leq Ce^{-\gamma_* t}\|(\rho_0, u_0)\|_{L^2}, \tag{14}$$

where $u^{\ell} = \widehat{u}(t, \xi)\mathbf{1}_{|\xi| \leq 1}$ and $u^{h} = \widehat{u}(t, \xi)\mathbf{1}_{|\xi| \geq 1}$.

How to obtain (12) rigorously?

First approach: Beauchard-Zuazua's method

Consider

$$\mathcal{L}_\xi(t) = |(\widehat{\rho}, \widehat{u})(\xi,t)|^2 + \frac{1}{2}\min\left(\frac{1}{|\xi|}, |\xi|\right) < \widehat{u} \cdot \widehat{\rho} >_{\mathbb{C}^n} . \qquad (15)$$

It gives the desired estimate, pointwise.

Second approach:

Homogeneous Littlewood-Paley decomposition

→ Allows to obtain precise decay rates, critical GWP results and to justify the strong relaxation limit.

Littlewood-Paley decomposition

Littlewood-Paley decomposition

- We define $\dot{\Delta}_j$ as dyadic blocks such that $f \in \mathcal{S}'_h(\mathbb{R}^d)$

$$f = \sum_{j\in\mathbb{Z}} \dot{\Delta}_j f \quad \text{and} \quad \mathrm{supp}(\widehat{\dot{\Delta}_j f}) \subset \{\xi \in \mathbb{R}^d \text{ t.q. } \frac{3}{4}2^j \leq |\xi| \leq \frac{8}{3}2^j\}.$$

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- The main motivation behind this decomposition is the following Bernstein inequality: $\forall k \in \mathbb{N}$, $p \in [1,\infty]$,

$$c2^{jk}\|\dot{\Delta}_j f\|_{L^p} \leq \|D^k \dot{\Delta}_j f\|_{L^p} \leq C2^{jk}\|\dot{\Delta}_j f\|_{L^p}.$$

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- The homogeneous Besov semi-norms are defined as follows:

$$\|f\|_{\dot{B}^s_{p,1}} \triangleq \sum_{j\in\mathbb{Z}} 2^{js}\|\dot{\Delta}_j f\|_{L^p}.$$

- We have $\dot{B}^0_{p,1} \subset L^p$, $\dot{B}^1_{2,1} \hookrightarrow \dot{H}^1$.

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- We have $\dot{B}^0_{p,1} \subset L^p$, $\dot{B}^1_{2,1} \hookrightarrow \dot{H}^1$.
- For a threshold $J_0 \in \mathbb{Z}$ and $s, s' \in \mathbb{R}$, we define:
- High-frequency norms: $\|f\|^h_{\dot{B}^s_{2,1}} \triangleq \sum_{j\geq J_0} 2^{js}\|\dot{\Delta}_j f\|_{L^2}$,
- Low-frequency norms: $\|f\|^\ell_{\dot{B}^{s'}_{p,1}} \triangleq \sum_{j\leq J_0} 2^{js'}\|\dot{\Delta}_j f\|_{L^p}$.

Toy-model analysis

Back to the damped *p*-system:

$$\begin{cases} \partial_t \rho + \partial_x u = 0, \\ \partial_t u + \partial_x \rho + u = 0.. \end{cases} \tag{16}$$

Applying the localisation operator $\dot{\Delta}_j$ to (16) and denoting $\dot{\Delta}_j f = f_j$, we have

$$\begin{cases} \partial_t \rho_j + \partial_x u_j = 0, \\ \partial_t u_j + \partial_x \rho_j + u_j = 0. \end{cases} \tag{17}$$

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Differentiating in time $\mathcal{L}_j(t) = \|(\rho_j, u_j, \partial_x \rho_j, \partial_x u_j)(t)\|_{L^2}^2 + \frac{1}{2}\int_{\mathbb{R}} u_j \partial_x \rho_j \, dx$, we get

$$\frac{d}{dt}\mathcal{L}_j(t) + \|(u_j, \partial_x u_j)\|_{L^2}^2 + \|\partial_x \rho_j\|_{L^2}^2 \leq 0. \tag{18}$$

Using Bernstein inequality, we have

$$\frac{d}{dt}\mathcal{L}_j(t) + \min(1, 2^{2j})\|(u_j, \rho_j)\|_{L^2}^2 \leq 0, \tag{19}$$

where $2^{2j} \sim |\xi|^2$.

We are going to use the following lemma.

Lemma

Let $p \geq 1$ *and* $X : [0, T] \to \mathbb{R}^+$ *be a continuous function such that* X^p *is a.e. differentiable. We assume that there exist a constant* $b \geq 0$ *and a measurable function* $A : [0, T] \to \mathbb{R}^+$ *such that*

$$\frac{1}{p}\frac{d}{dt}X^p + bX^p \leq AX^{p-1} \quad a.e.\ on \ \ [0, T].$$

Then, for all $t \in [0, T]$, *we have*

$$X(t) + b\int_0^t X \leq X_0 + \int_0^t A.$$

Applying this lemma to

$$\frac{d}{dt}\mathcal{L}_j(t) + \min(1, 2^{2j})\|(u_j, \rho_j)\|_{L^2}^2 \leq 0, \tag{20}$$

since $\mathcal{L}_j \sim \|(u_j, \rho_j)\|_{L^2}^2$, we obtain

$$\sqrt{\mathcal{L}_j(t)} + \min(1, 2^{2j}) \int_0^t \|(u_j, \rho_j)\|_{L^2} \leq 0. \tag{21}$$

Using that $\sqrt{\mathcal{L}_j(t)} \sim \|(u_j, \rho_j)\|_{L^2}$, we get

$$\|(u_j, \rho_j)(t)\|_{L^2} + \min(1, 2^{2j}) \int_0^t \|(u_j, \rho_j)\|_{L^2} \leq 0. \tag{22}$$

From here, we distinguish two cases.

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From here, we distinguish two cases.

- For high frequencies: $j \geq 0 \implies \min(1,2^{2j}) = 1$.

Multiplying (22) by 2^{js} for $s \in \mathbb{R}$ and summing on $j \geq 0$, we obtain

$$\|(u,\rho)(t)\|^h_{\dot{B}^s_{2,1}} + \|(u,\rho)\|^h_{L^1_T(\dot{B}^s_{2,1})} \leq 0.$$

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$$\|(u, \rho)(t)\|^\ell_{\dot{B}^s_{2,1}} + \|(u, \rho)\|^\ell_{L^1_T(\dot{B}^{s+2}_{2,1})} \leq 0.$$

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$$\|(u_j,\rho_j)(t)\|_{L^2} + \min(1,2^{2j})\int_0^t \|(u_j,\rho_j)\|_{L^2} \lesssim 0. \tag{22}$$

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• Heat effect in low frequencies (not very good) and exponential stabilization in high frequencies.

• From here: optimal decay rates using time-weights and interpolations.

• Notice the $L^1_T(B^{s+2}_{2,1})$ norm compared to the usual $L^2_T(H^{s+1})$ norm.

Back to the general system

$$\partial_t U + A\partial_x U + BU = 0.$$

Under the Kalman rank condition (or the Shizuta-Kawashima) condition for (A, B), differentiating in time the following functional

$$\mathcal{L}_j(t) = \|U_j(t)\|_{H^1}^2 + \eta \int_{\mathbb{R}} \left(\sum_{k=1}^{n-1} < BA^{k-1}U_j, BA^k U_j > \right)$$

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Thus

$$\|U(t)\|_{\dot{B}^s_{2,1}}^h + \|U\|_{L^1_T(\dot{B}^s_{2,1})}^h \leq 0,$$

and

$$\|U(t)\|_{\dot{B}^s_{2,1}}^\ell + \|U\|_{L^1_T(\dot{B}^{s+2}_{2,1})}^\ell \leq 0.$$

- What we have just seen allows us to recover the classical existence results for nonlinear systems in a slightly better framework:

$$\dot{B}^{\frac{d}{2}}_{2,1} \cap \dot{B}^{\frac{d}{2}+1}_{2,1} \quad \mathrm{vs} \quad H^s \quad \text{for } s > \frac{d}{2} + 1.$$

- Recalling that

$$H^s(s > \frac{d}{2} + 1) \hookrightarrow B^{\frac{d}{2}+1}_{2,1} \hookrightarrow \dot{B}^{\frac{d}{2}}_{2,1} \cap \dot{B}^{\frac{d}{2}+1}_{2,1} \hookrightarrow \dot{B}^{\frac{d}{p},\frac{d}{2}+1}_{p,2}(p > 2) \hookrightarrow \mathcal{C}^1_b.$$

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Essentially:

- We have to go beyond "standard hypocoercivity" in the low frequencies.
- The eigenvalues in low-frequency are purely real → it is possible to decouple the system, up to linear high-order terms (good in LF).
- For that matter we introduce a purely damped mode, in contrast with the heat behavior, in the low-frequency regime,

Low-frequency analysis.

Low frequencies in a simple case

Back to the localized damped p-system:

$$\begin{cases} \partial_t u_j + \partial_x v_j = 0 \\ \partial_t v_j + \partial_x u_j + v_j = 0, \end{cases}$$

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$$\begin{cases} \partial_t u_j - \partial^2_{xx} u_j = -\partial_x w_j \\ \partial_t w_j + w_j = -\partial^2_{xx} w_j - \partial^3_{xxx} \rho_j. \end{cases}$$

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• This diagonalisation of the system directly exhibit the low-frequency behaviour observed in the spectral analysis.

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• To deal with the linear source terms, we use the Bernstein inequality

$$\|\partial_x f\|^{\ell}_{B^s_{p,1}} = \|f\|^{\ell}_{B^{s+1}_{p,1}} = \sum_{j \leq J_0} 2^{j(s+1)} \|f_j\|_{L^p} \leq \sum_{j \leq J_0} 2^{js} 2^j \|f_j\|_{L^p} \leq J_0 \|f\|^{\ell}_{B^s_{p,1}}.$$

where J_0 is the threshold between low and high frequencies that has to be chosen small enough.

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Defining the damped mode $w_j = v_j + \partial_x u_j$, the system can be rewritten

$$\begin{cases} \partial_t u_j - \partial_{xx}^2 u_j = -\partial_x w_j \\ \partial_t w_j + w_j = -\partial_{xx}^2 w_j - \partial_{xxx}^3 \rho_j. \end{cases}$$

• This diagonalisation of the system directly exhibit the low-frequency behaviour observed in the spectral analysis.

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$$\|\partial_x f\|_{B^s_{p,1}}^{\ell} = \|f\|_{B^{s+1}_{p,1}}^{\ell} = \sum_{j\leq J_0} 2^{j(s+1)}\|f_j\|_{L^p} \leq \sum_{j\leq J_0} 2^{js}2^{j}\|f_j\|_{L^p} \leq J_0\|f\|_{B^s_{p,1}}^{\ell}.$$

where J_0 is the threshold between low and high frequencies that has to be chosen small enough.

• A priori estimates in a L^p framework for $2 \leq p \leq 4$ is available in the low-frequency regime.

General case

General case

In the general case, the system can be rewritten as follows:

$$\begin{cases} \partial_t U_1 + A_{1,1}\partial_x U_1 + A_{1,2}\partial_x U_2 = 0, \\ \partial_t U_2 + A_{2,1}\partial_x U_1 + A_{2,2}\partial_x U_2 + DU_2 = 0. \end{cases}$$

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$$W \triangleq U_2 + D^{-1}A_{2,1}\partial_x U_1 + D^{-1}A_{2,2}\partial_x U_2 = D^{-1}\partial_t U_2.$$

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The system can be rewritten

$$\begin{cases} \partial_t U_1 - A_{1,2}D^{-1}A_{2,1}\partial_x\partial_x U_1 = f \\ \partial_t W + DW = g \end{cases} \tag{23}$$

where f and g are controllable in the low-frequency regime with Bernstein-type inequalities.

What can we say about the operator $A_{1,2}D^{-1}A_{2,1}\partial_x\partial_x$ in the equation of U_1?

General case

To study the equation of U_1, we have the following property

Lemma

For $D > 0$, the following assertions are equivalent:

- *(A,B) satisfy the Kalman rank condition,*
- *the operator $\mathcal{A} := A_{1,2}D^{-1}A_{2,1}\partial_{xx}^2$ is strongly elliptic.*

→ We may study the equations of W and U_1 separately, the former as a damped equation and the latter as a heat equation.

- This approach can be applied to general systems of the form:

$$\begin{cases} \partial_t U + \sum_{j=1}^{d} A^j(U)\partial_{x_j} U + G(U) = 0, \\ U_0(x,t) = U_0(x), \end{cases} \tag{24}$$

for solutions close to a constant equilibrium $\bar{U}$ such that $G(\bar{U}) = 0$.

Important assumptions:

- $A_{1,1}(\bar{U}) = 0$ which means that $\bar{u} = 0$ for fluid-type systems.
- On the other hand, we need $\bar{\rho} > 0$.

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Tools to deal with the nonlinear terms:

- Embeddings for the type:

$$\dot{B}^{\frac{d}{p}}_{p,1} \hookrightarrow L^\infty, \quad \dot{B}^{\frac{d}{p}+1}_{p,1} \hookrightarrow \dot{W}^{1,\infty} \quad \text{and} \quad B^s_{2,1} \hookrightarrow B^s_{p,1}$$

- Advanced product laws, commutators estimates and composition estimates to deal with the $(L^2)^h \cap (L^p)^\ell$ setting:

$$\|ab\|^h_{\dot{B}^s_{2,1}} \lesssim \|a\|_{\dot{B}^{\frac{d}{p}}_{p,1}} \|b\|^h_{\dot{B}^s_{2,1}} + \|b\|_{\dot{B}^{\frac{d}{p}}_{p,1}} \|a\|^h_{\dot{B}^s_{2,1}} + \|a\|^\ell_{\dot{B}^{\frac{d}{p}-\frac{d}{p*}}_{p,1}} \|b\|^\ell_{\dot{B}^\sigma_{p,1}} + \|b\|^\ell_{\dot{B}^{\frac{d}{p}-\frac{d}{p*}}_{p,1}} \|a\|^\ell_{\dot{B}^\sigma_{p,1}}.$$

Well-posedness result for nonlinear systems.

We set $Z = U - \bar{U}$.

Theorem (Danchin, C-B '22 Math. Ann.)

Let $d \geq 1$, $p \in [2,4]$. *There exists* $c_0 = c_0(p) > 0$ *and* J_0 *such that if*

$$\|Z_0\|^{\ell}_{\dot{B}^{\frac{d}{p}}_{p,1}} + \|Z_0\|^{h}_{\dot{B}^{\frac{d}{2}+1}_{2,1}} \leq c_0,$$

then the system admits a unique solution Z *satisfying*

$$X_p(t) \lesssim \|Z_0\|^{\ell}_{\dot{B}^{\frac{d}{p}}_{p,1}} + \|Z_0\|^{h}_{\dot{B}^{\frac{d}{2}+1}_{2,1}} \quad \text{for all } t \geq 0,$$

where

$$\begin{aligned} X_p(t) \triangleq\ & \|Z\|^{h}_{L^\infty_t(\dot{B}^{\frac{d}{2}+1}_{2,1})} + \|Z\|^{h}_{L^1_t(\dot{B}^{\frac{d}{2}+1}_{2,1})} + \|Z_2\|_{L^2_t(\dot{B}^{\frac{d}{p}}_{p,1})} \\ & + \|Z\|^{\ell}_{L^\infty_t(\dot{B}^{\frac{d}{p}}_{p,1})} + \|Z_1\|^{\ell}_{L^1_t(\dot{B}^{\frac{d}{p}+2}_{p,1})} + \|Z_2\|^{\ell}_{L^1_t(\dot{B}^{\frac{d}{p}+1}_{p,1})} + \|W\|_{L^1_t(\dot{B}^{\frac{d}{p}}_{p,1})}. \end{aligned}$$

Proof: Previous linear analysis + Perturbation and Bootstrap arguments.

- The hypocoercive-type analysis can be extended to general system of any order

$$\partial_t V + A(D)V + L(D)V = 0, \quad \text{where}$$

  - $A(D)$ is a skew-symmetric homogeneous Fourier multiplier of order α,
  - $L(D)$ is a partially elliptic homogeneous Fourier multiplier of order β.
- **What dictates the decay rates is difference of order between A and B.**

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- Anisotropic case (cf. Bianchini-CB-Paicu) concerning stably stratified solutions of the 2D-Boussinesq system.
- **Open question:** What kind of nonlinearities can we include? Relation between partial dissipation, hyperbolicity and anisotropy.

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- Another interesting case

$$\partial_t U + A\partial_x U + BU = 0$$

  for A symmetric and B non-symmetric e.g. Euler-Maxwell system or Timoshenko system
- One must consider Kalman rank condition for (B^s, B^a) where B^s is the symmetric part of B and B^a the skew-symmetric part.

Second part: Relaxation procedure and hyperbolisation

Cattaneo approximation of the heat equation

Let us consider the heat equation on $\mathbb{R}^d$

$$\partial_t \rho - \Delta \rho = 0.$$

Its Cattaneo hyperbolic approximation reads

$$\begin{cases} \partial_t \rho_\varepsilon + \partial_x u_\varepsilon = 0, \\ \varepsilon^2 \partial_t u_\varepsilon + \partial_x \rho_\varepsilon + u_\varepsilon = 0. \end{cases} \tag{25}$$

When $\varepsilon \to 0$, we recover a heat equation for ρ and a Darcy-type law $u = \partial_x \rho$.

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- System (25) has a partially dissipative and hyperbolic structure.
- → *Dissipative hyperbolisation*.
- **How to justify the limit $\varepsilon \to 0$ rigorously?**
- Tools from the previous part!

Solution First! Spectral analysis

Cattaneo approximation:

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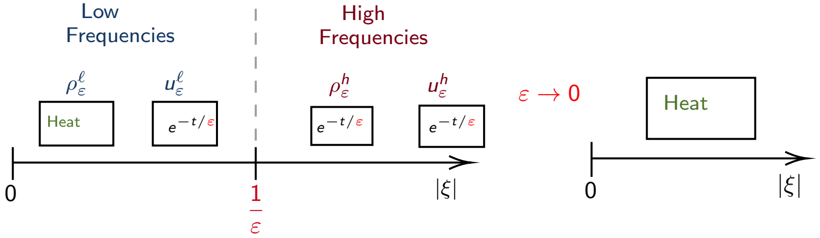

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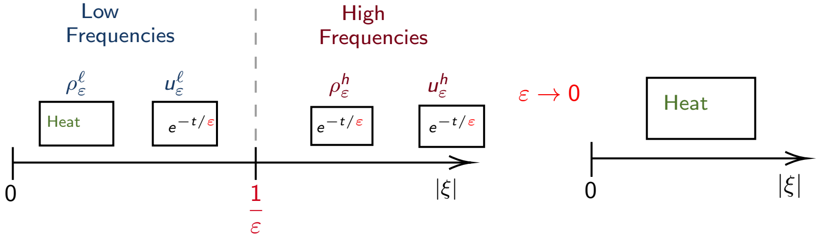

- The Cattaneo approximation creates a high-frequency regime where the solution is exponentially damped.
- **Goal:** Justify this process for nonlinear systems.

Spaces

- We work with the following hybrid homogeneous Besov norms:

$$\|f\|^{h}_{\dot{B}^{s}_{2,1}} \triangleq \sum_{j\geq \frac{\eta}{\varepsilon}} 2^{js}\|\dot{\Delta}_j f\|_{L^2} \quad \text{and} \quad \|f\|^{\ell}_{\dot{B}^{s'}_{p,1}} \triangleq \sum_{j\leq \frac{\eta}{\varepsilon}} 2^{js'}\|\dot{\Delta}_j f\|_{L^p}.$$

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- For low-frequencies: $j \leq \dfrac{\eta}{\varepsilon}$,

$$\begin{cases}\partial_t\rho_j + \partial_x u_j = 0\\ \varepsilon^2\partial_t u_j + \partial_x\rho_j + u_j = 0,\end{cases}$$

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  Careful, the Bernstein inequality is different:

$$\|\partial_x f\|^{\ell}_{B^s_{p,1}} \leq \frac{\eta}{\varepsilon}\|f\|^{\ell}_{B^s_{p,1}}.$$

Details of computations

For $s \in \mathbb{R}$, we have

$$\|(u,\varepsilon w)(t)\|^{\ell}_{B^s_{p,1}} + \|\rho\|^{\ell}_{L^1_T(B^{s+2}_{p,1})} + \frac{1}{\varepsilon}\|w\|^{\ell}_{L^1_T(B^s_{p,1})} \leq \|(u_0,w_0)\|^{\ell}_{B^s_{p,1}} + \varepsilon\|w\|^{\ell}_{L^1_T(B^{s+2}_{p,1})} + \varepsilon\|\rho\|^{\ell}_{L^1_T(B^{s+3}_{p,1})}$$

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With the Berstein inequality, we have

$$\varepsilon\|\rho\|^{\ell}_{L^1_T(B^{s+3}_{p,1})} \leq \eta\|\rho\|^{\ell}_{L^1_T(B^{s+2}_{p,1})} \quad \text{and} \quad \varepsilon\|w\|^{\ell}_{L^1_T(B^{s+2}_{p,1})} \leq \frac{\eta^2}{\varepsilon}\|w\|^{\ell}_{L^1_T(B^s_{p,1})}.$$

Thus, choosing η small enough, these terms can be absorbed by the l.h.s.

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- This estimate provides $\mathcal{O}(\varepsilon)$ bounds on $w = u + \partial_x\rho$ which is crucial to justify the relaxation.

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- **High frequencies $j \geq \frac{\eta}{\varepsilon}$:** Hypocoercivity-type approach but there is no damped mode!

High frequencies trick

To be able to recover $\mathcal{O}(\varepsilon)$ bounds on w in high frequencies, we use the Bernstein inequality

$$\|f\|^{h}_{B^{s}_{2,1}} \leq \frac{\varepsilon}{\eta}\|\partial_x f\|^{h}_{B^{s}_{2,1}}.$$

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$$\|f\|^h_{B^s_{2,1}} \leq \frac{\varepsilon}{\eta}\|\partial_x f\|^h_{B^s_{2,1}}.$$

Say you want to obtain uniform bounds for w in $B^{\frac{d}{2}}_{2,1}$, then you should assume that the initial data are in $B^{\frac{d}{2}+1}_{2,1}$ and use that

$$\|w\|^h_{B^{\frac{d}{2}}_{2,1}} \leq \frac{\varepsilon}{\eta}\|w\|^h_{B^{\frac{d}{2}+1}_{2,1}}.$$

$\Longrightarrow$ We must study the low and high frequencies are at different regularities.

Back to the compressible Euler equations

Back to the compressible Euler equations

The system reads:

$$\begin{cases} \partial_t \rho + \mathrm{div}\,(\rho u) = 0, \\ \varepsilon^2(\partial_t u + u \cdot \nabla u) + \dfrac{\nabla P(\rho)}{\rho} + u = 0. \end{cases} \tag{E}$$

The damped mode associated to the relaxation is $w = u + \dfrac{\nabla P(\rho)}{\rho}$.

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Inserting it in the above equation, we recover

$$\partial_t \rho - \Delta P(\rho) = \mathrm{div}\, w.$$

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- Let $\mathcal{N}$ be the solution of the porous media equation:

$$\partial_t\mathcal{N} - \Delta P(\mathcal{N}) = 0.$$

Then, using that $\|w\|_{L^1_T(B^s_{p,1})} = \mathcal{O}(\varepsilon)$, in the error estimates for $\widetilde{\rho} = \rho - \mathcal{N}$, we can justify that ρ converges strongly toward $\mathcal{N}$ in $B^{s-1}_{p,1}$.

Relaxation result

Theorem (Danchin, C-B, Math. Ann. 2022)

Let $d \geq 1$, $p \in [2,4]$ and $\varepsilon > 0$.

- *Let $\bar{\rho}$ be a strictly positive constant and $(\rho^\varepsilon - \bar{\rho}, u^\varepsilon)$ be the solution of the compressible Euler system with damping (constructed with the previous arguments)*
- *Let $\mathcal{N} \in \mathcal{C}_b(\mathbb{R}^+; \dot{B}^{\frac{d}{p}}_{p,1}) \cap L^1(\mathbb{R}^+; \dot{B}^{\frac{d}{p}+2}_{p,1})$ be the unique solution associated to the Cauchy problem:*

$$\begin{cases} \partial_t \mathcal{N} - \Delta P(\mathcal{N}) = 0 \\ \mathcal{N}(0,x) = \mathcal{N}_0 \in \dot{B}^{\frac{d}{p}}_{p,1} \end{cases}$$

If we assume that

$$\|\rho_0^\varepsilon - \mathcal{N}_0\|_{B^{\frac{d}{p}-1}_{p,1}} \leq C\varepsilon,$$

then

$$\|\rho^\varepsilon - \mathcal{N}\|_{L^\infty(\mathbb{R}_+; \dot{B}^{\frac{d}{p}-1}_{p,1})} + \|\rho^\varepsilon - \mathcal{N}\|_{L^1(\mathbb{R}_+; \dot{B}^{\frac{d}{p}+1}_{p,1})} + \left\| \frac{\nabla P(\rho^\varepsilon)}{\rho^\varepsilon} + u^\varepsilon \right\|_{L^1(\mathbb{R}^+; \dot{B}^{\frac{d}{p}}_{p,1})} \leq C\varepsilon.$$

Remarks

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- Performing a similar analysis with Sobolev spaces does not allow (to the best of my knowledge) to exhibit an explicit convergence rate.
- It only leads to $\|w\|_{L^2_T(H^s)} = \mathcal{O}(1)$ **vs** $\|w\|_{L^1_T(B^s_{2,1})} = \mathcal{O}(\varepsilon)$

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- It only leads to $\|w\|_{L^2_T(H^s)} = \mathcal{O}(1)$ **vs** $\|w\|_{L^1_T(B^s_{2,1})} = \mathcal{O}(\varepsilon)$
- First result to establish the strong relaxation limit in the multi-dimensional setting.
- It can be employed in many other contexts.

Application to a (partially) hyperbolic Navier-Stokes system

Hyperbolic Navier-Stokes equations

We have just seen that the equation

$$\partial_t u - \Delta u = 0$$

can be approximated, for a small ε, by the following hyperbolic system

$$\begin{cases} \partial_t u + \operatorname{div} v = 0 \\ \varepsilon^2 \partial_t v + \nabla u + v = 0. \end{cases}$$

- Our aim is now to understand to what extent this approximation can be used to approximate systems modelling physical phenomena.

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- Our aim is now to understand to what extent this approximation can be used to approximate systems modelling physical phenomena.

Performing such approximation for the compressible Navier-Stokes system, one has

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho u) = 0, \\ \partial_t(\rho u) + \operatorname{div}(\rho u \otimes u) + \nabla p = \operatorname{div} \tau, \\ \partial_t(\rho T) + \operatorname{div}(\rho u T + up) + \operatorname{div} q - \operatorname{div}(\tau \cdot u) = 0, \\ \varepsilon^2 \partial_t q + q + \kappa \nabla T = 0, \end{cases} \tag{NSCC}$$

Let us now see how to justify that the solution of this system converge to the solution of the classical Navier-Stokes equations.

Frequency-splitting

- For the limit system: Danchin showed the existence of solutions satisfying different properties for $|\xi| \leq K$ and $|\xi| \geq K$ where K is a large constant.

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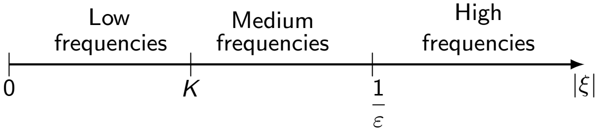
In order to obtain the complete picture, we divide the frequency space as:



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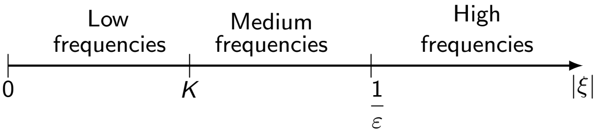

Formally, when $\varepsilon \to 0$, it means that:

- The low frequency regime is not modified.
- The mid-frequency regime becomes larger and larger and recovers the high-frequency regime.
- The high frequency regime disappears in we recover the limit system.

Tools

Tools

- We define homogeneous Besov spaces restricted in frequency as follows:

$$\|f\|^{\ell}_{\dot{B}^s_{p,1}} := \sum_{j\leq J_0} 2^{js}\|f_j\|_{L^2}, \qquad \|f\|^{m,\varepsilon}_{\dot{B}^s_{p,1}} := \sum_{J_0\leq j\leq J_\varepsilon} 2^{js}\|f_j\|_{L^2}$$

$$\text{and} \quad \|f\|^{h,\varepsilon}_{\dot{B}^s_{p,1}} := \sum_{j\geq J_\varepsilon-1} 2^{js}\|f_j\|_{L^2}$$

  where $J_0 = K$, for $K > 0$ a constant, and $J_\varepsilon = -\dfrac{1}{\varepsilon}$.
- In each regime, different methods have to be developed and patched together to derive a priori estimates.
- Hypocoercivity, efficient unknowns and limit system's analysis.
- Difficulty: handling the nonlinearities.

The Jin-Xin Approximation.

Jin-Xin Approximation

In collaboration with Ling-yun Shou (JDE), we justified the strong convergence of the diffusive Jin-Xin approximation

$$\begin{cases} \dfrac{\partial}{\partial t} u + \displaystyle\sum_{i=1}^{d} \dfrac{\partial}{\partial x_i} v_i = 0, \\ \varepsilon^2 \dfrac{\partial}{\partial t} v_i + A_i \dfrac{\partial}{\partial x_i} u = -\big(v_i - f_i(u)\big), \qquad i = 1, 2, ..., d, \end{cases} \tag{26}$$

toward viscous conservation laws:

$$\frac{\partial}{\partial t} u^* + \sum_{i=1}^{d} \frac{\partial}{\partial x_i} f_i(u^*) = \sum_{i=1}^{d} \frac{\partial}{\partial x_i} (A_i \frac{\partial}{\partial x_i} u^*), \tag{27}$$

Other applications:

- 2D-Boussinesq System (Bianchini-CB-Paicu)
- Baer-Nunziato System (Burtea-CB-Tan), M3AS.
- Chemotaxis/Keller-Segel, (CB-He-Shou) SIAM.

Conclusion

- Hypocoercivity tells you that when the dissipation is not strong enough, its interactions with the hyperbolic part can compensate the lack of coercivity.
- When the skew-symmetric operator A and the dissipative B are of different order then the decay rates may not be exponential and the rats depend on the difference of their order.
- In the full space $\mathbb{R}$, the classical hypocoercivity techniques need to be extended to treat the low frequencies.
- The hyperbolic relaxation creates a temporary exponentially stable regime and the low frequencies correspond to the behavior of the limit system.

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Overall, splitting the frequency space is nice!

Thank you for your attention.

Formal link between (IPM) and (2D-B)

The 2-dimensional Boussinesq system read

$$\begin{cases} \partial_t \eta + \mathbf{u} \cdot \nabla \eta = 0, \\ \partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} + \nabla P = \eta \mathbf{g}, \qquad \mathbf{g} = (0, -g), \\ \nabla \cdot \mathbf{u} = 0. \end{cases} \tag{E}$$

The linearized system around $\overline{\rho}_{\text{eq}}(y) = \rho_0 - y$, reads

$$\begin{cases} \partial_t b - \mathcal{R}_1 \Omega = 0, \\ \varepsilon^2 \partial_t \Omega - \mathcal{R}_1 b + \Omega = 0. \end{cases} \tag{28}$$

where

$$\mathcal{R}_1 = \frac{\partial_x}{(-\Delta)^{-\frac{1}{2}}}$$

Formally, as $\varepsilon \to 0$, the second equation gives the Darcy's law $\widetilde{\Omega}^\varepsilon = \mathcal{R}_1 \widetilde{b}^\varepsilon$ and inserting it in the first one gives the linear part of the incompressible porous media equation:

$$\partial_t \widetilde{b}^\varepsilon - \mathcal{R}_1^2 \widetilde{b}^\varepsilon = 0.$$

The HPC System

In joint work with Q. He and L-Y. Shou, we studied the following hyperbolic-parabolic system:

$$\begin{cases} \partial_t \rho + \operatorname{div}(\rho u) = 0, \\ \partial_t(\rho u) + \operatorname{div}(\rho u \otimes u) + \nabla P(\rho) + \dfrac{1}{\varepsilon}\rho u - \mu\rho\nabla\phi = 0, \\ \partial_t \phi - \Delta\phi - a\rho + b\phi = 0, \qquad\qquad x \in \mathbb{R}^d, \quad t > 0, \end{cases} \tag{HPC}$$

In this case, when $\varepsilon \to 0$, we show that the diffusive-rescaled solution of (HPC) converges strongly to the solution of the Keller-Segel system:

$$\begin{cases} \partial_t \rho - \operatorname{div}(\nabla P(\rho) - \mu\rho\nabla\phi) = 0, \\ \rho u = -\nabla P(\rho) + \mu\rho\nabla\phi, \\ -\Delta\phi - a\rho + b\phi = 0, \end{cases} \tag{KS}$$

Multifluid system

In a joint work with C. Burtea, J. Tan and L.-Y. Shou, we studied the following damped Baer-Nunziato system:

$$\begin{cases} \partial_t \alpha_\pm + u \cdot \nabla \alpha_\pm = \pm \dfrac{\alpha_+\alpha_-}{2\mu + \lambda}\left(P_+(\rho_+) - P_-(\rho_-)\right), \\ \partial_t(\alpha_\pm \rho_\pm) + \operatorname{div}(\alpha_\pm \rho_\pm u) = 0, \\ \partial_t(\rho u) + \operatorname{div}(\rho u \otimes u) + \nabla P + \eta \rho u = 0, \\ \rho = \alpha_+\rho_+ + \alpha_-\rho_-, \\ P = \alpha_+ P_+(\rho_+) + \alpha_- P_-(\rho_-) \end{cases} \tag{BN}$$

Limit $\lambda, \mu, \nu \to 0$.

- Difficulties: the entropy that is naturally associated with this system is only positive semi-definite.
- The system (BN) is not a system of conservation laws
- We find an ad-hoc change of variables that enables us to symmetrize the system with a good structure to treat the nonlinear terms.

Overdamping

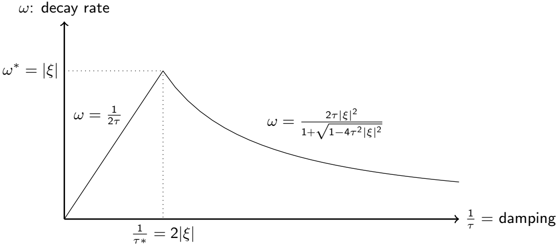

Figure: A graph of overdamping phenomenon for System (**??**).

Decay estimates

Theorem (Danchin, C-B '22)

Assuming additionally that $Z_0 \in \dot{B}^{-\sigma_1}_{2,\infty}$ *for* $\sigma_1 \in \left]-\frac{d}{2}, \frac{d}{2}\right]$ *then there exists* $C > 0$ *such that*

$$\|Z(t)\|_{\dot{B}^{-\sigma_1}_{2,\infty}} \leq C \, \|Z_0\|_{\dot{B}^{-\sigma_1}_{2,\infty}}, \quad \forall t \geq 0.$$

Moreover, if $\sigma_1 > 1 - d/2$,

$$\langle t \rangle \triangleq \sqrt{1+t^2}, \quad \alpha_1 \triangleq \frac{\sigma_1 + \frac{d}{2} - 1}{2} \quad \text{and} \quad C_0 \triangleq \|Z_0\|^{\ell}_{\dot{B}^{-\sigma_1}_{2,\infty}} + \|Z_0\|^{h}_{\dot{B}^{\frac{d}{2}+1}_{2,1}},$$

then Z *satisfies the following decay estimates:*

$$\begin{aligned} \sup_{t\geq 0} \left\| \langle t \rangle^{\frac{\sigma+\sigma_1}{2}} Z(t) \right\|^{\ell}_{\dot{B}^{\sigma}_{2,1}} &\leq CC_0 \quad \text{if} \quad -\sigma_1 < \sigma \leq d/2 - 1, \\ \sup_{t\geq 0} \left\| \langle t \rangle^{\frac{\sigma+\sigma_1}{2}+\frac{1}{2}} Z_2(t) \right\|^{\ell}_{\dot{B}^{\sigma}_{2,1}} &\leq CC_0 \quad \text{if} \quad -\sigma_1 < \sigma \leq d/2 - 2, \\ \text{and} \quad \sup_{t\geq 0} \left\| \langle t \rangle^{2\alpha_1} Z(t) \right\|^{h}_{\dot{B}^{\frac{d}{2}+1}_{2,1}} &\leq CC_0. \end{aligned}$$